

HEUSLER COMPOUNDS (AND RELATED) IN MAGNETIC TUNNEL JUNCTIONS

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DFG
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Programme
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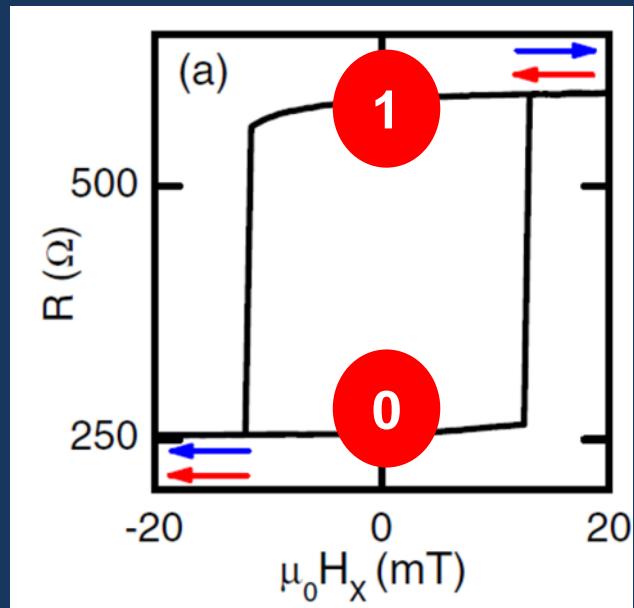
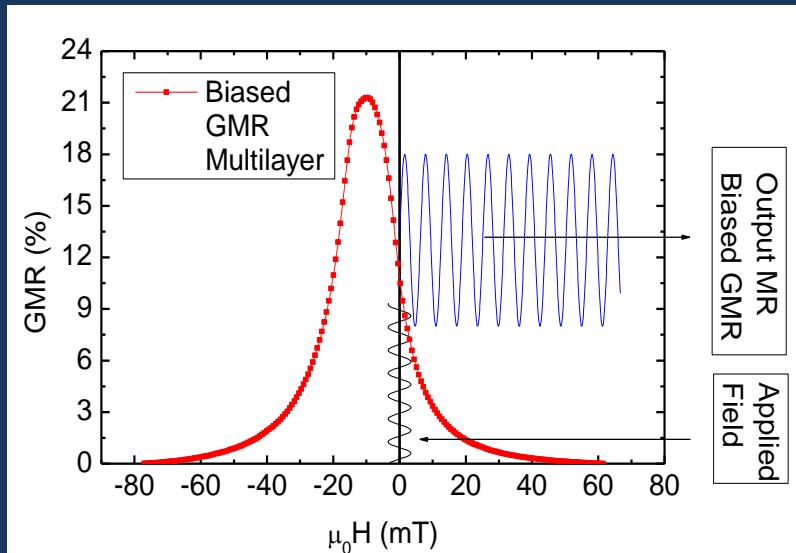


FAQ's: Where is Bielefeld ?

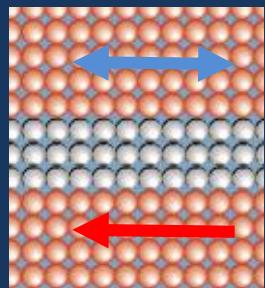
There



Enhanced MR is very general in very thin (1nm) layered structures:



Metallic spacer
(Cu, Cr, ..):
Giant MagnetoResistance



Insulating spacer
(MgO, Al₂O₃, ..):
Tunneling MagnetoResistance

Now, we have a RAM



The reference– ultrathin CoFeB

The samples:



MTJ stack sequence for this study:
pseudo-spinvalves with
ultrathin CoFeB and varying MgO

**Magnetic Tunnel Junctions
with ultrathin CoFeB are
perpendicular**

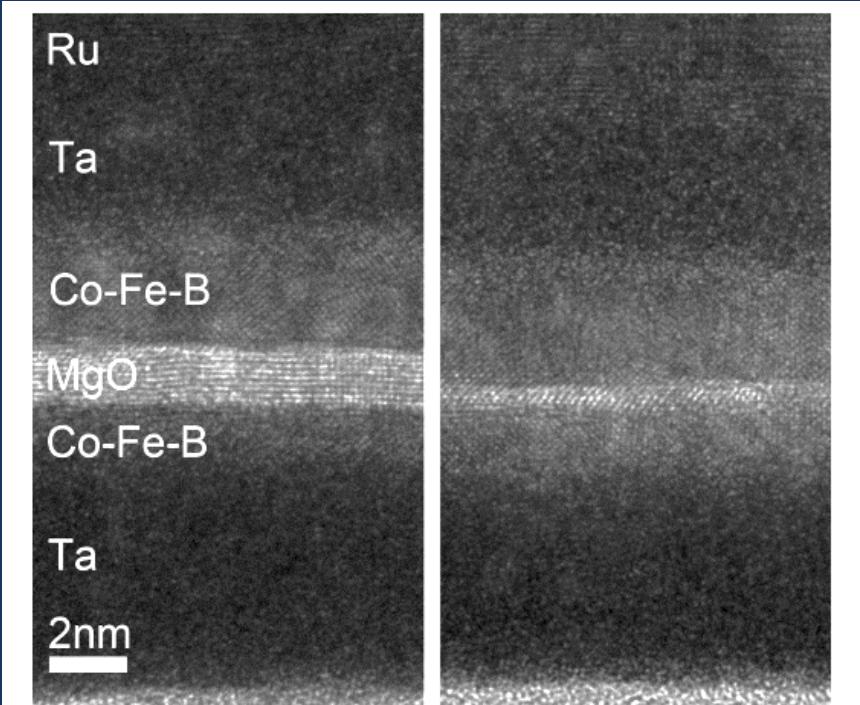
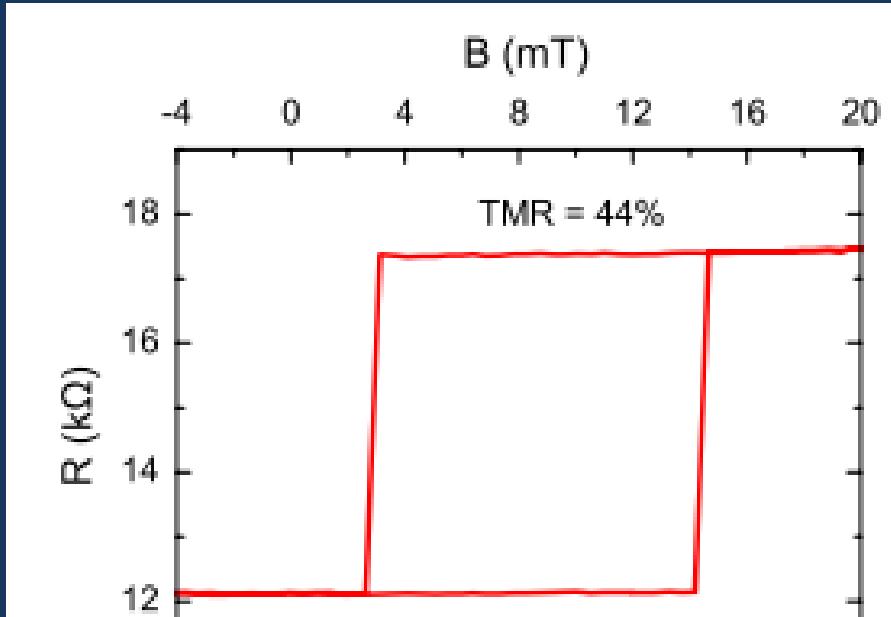


Fig. 1. HRTEM images of a thick 10 ML (left) and a heated 3 ML MgO barrier (right). The IQR values are $(5.6 \pm 1.5)^\circ$ (10 ML) and $(6.7 \pm 0.8)^\circ$ (3 ML).

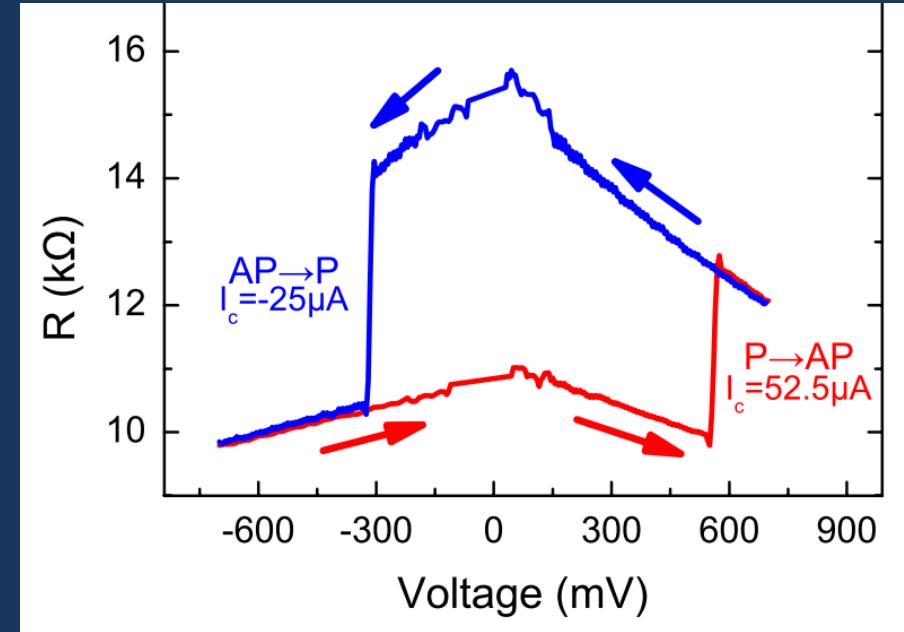
**and work down to MgO barrier-
thickness of 3 monolayers**

The reference– ultrathin CoFeB STT switching

Results for STT-switching this ultrathin CoFeB/MgO/CoFeB system for low-RA MgO and small junctions:



Resistance vs. external magnetic field for perpendicular MTJs
1.0nm CoFeB / 4 ML MgO / 1.2nm CoFeB
gives around 40-50% TMR



RV-characteristic with an applied field of 8.6 mT
- average critical current density:

$$\underline{\underline{2 \cdot 10^5 \text{ A/cm}^2 (!!)}}$$

Now we have a working and stable STT-MRAM



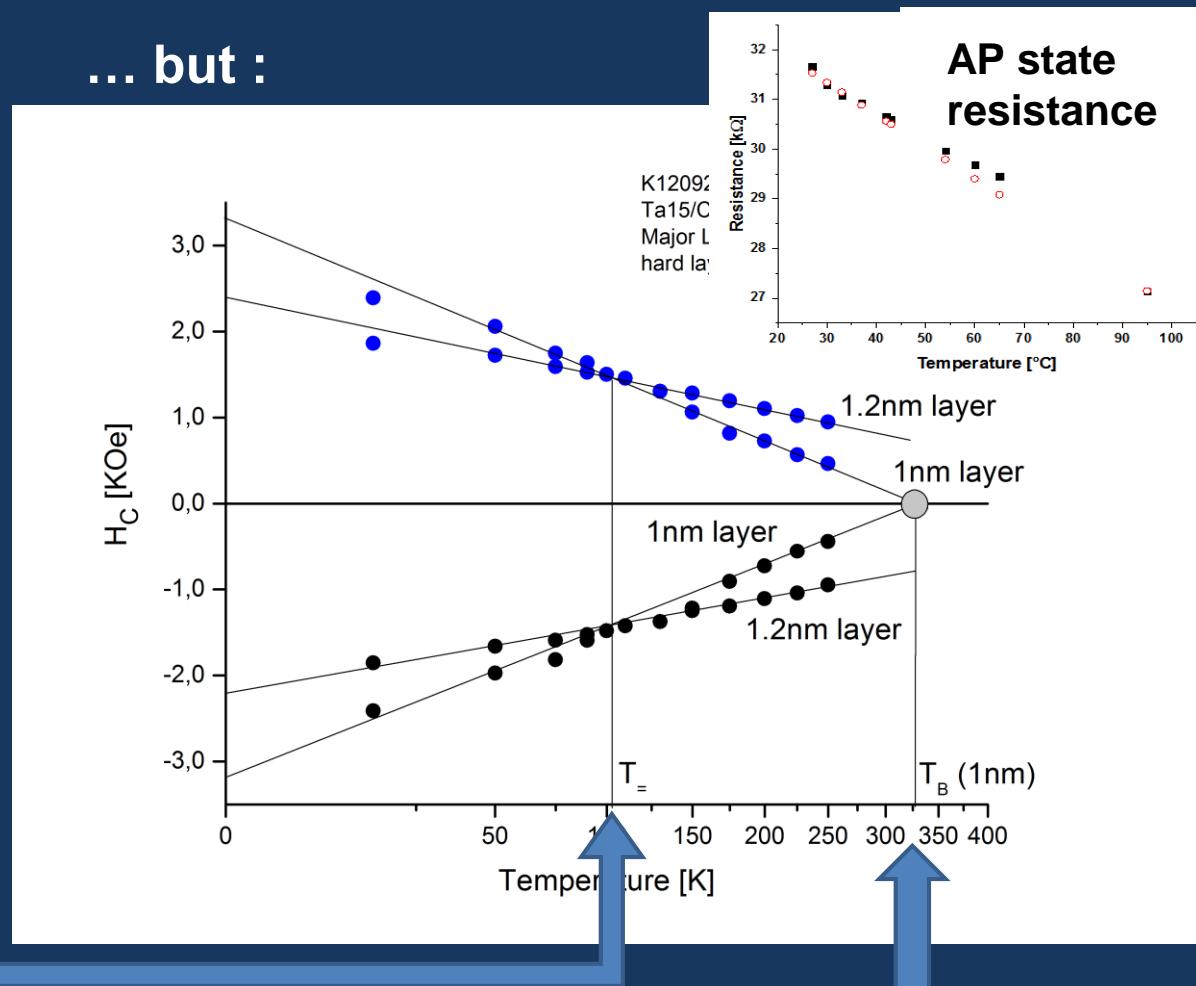
Collaboration with
M. Münzenberg group



1. the layers change their role ..
low T: thick layer is the free layer
high T: thin layer is free

The reference – ultrathin CoFeB magnetic properties

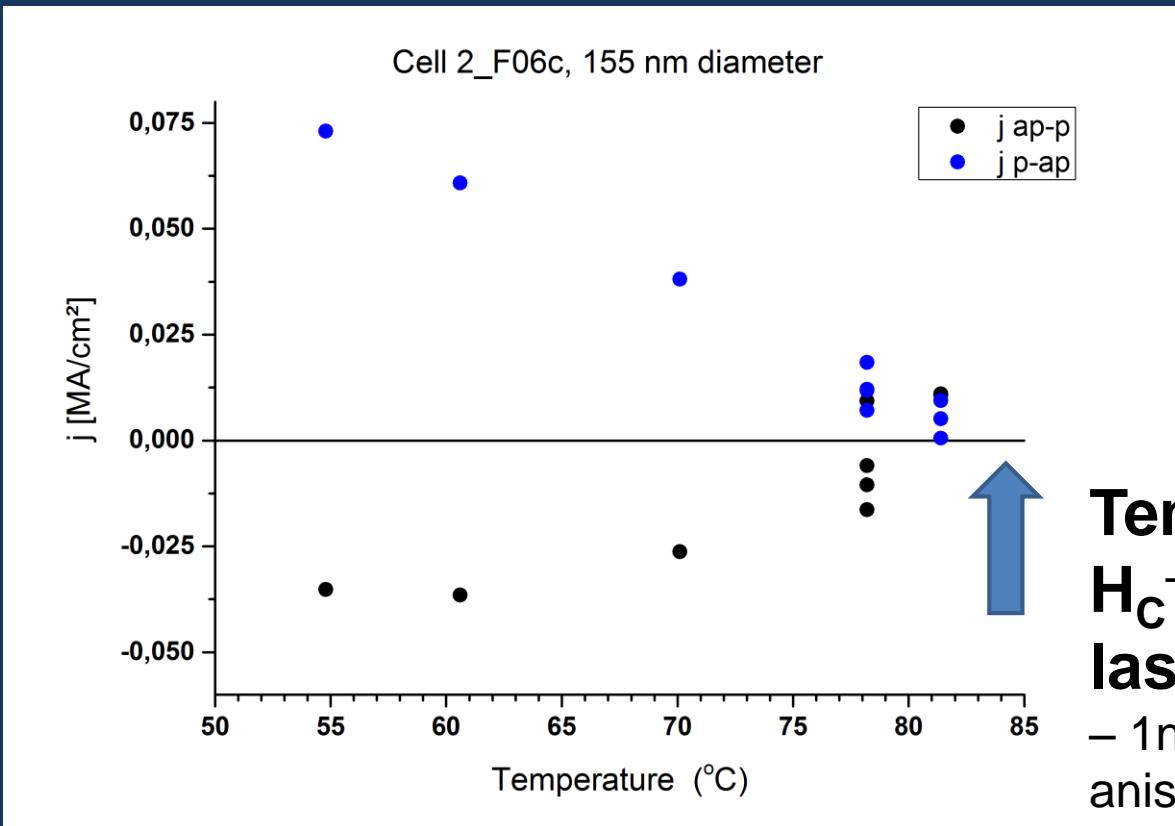
... but :



2. the thin layer shows $H_C^\perp = 0$? anisotropy switches to ip or? superparamagnetic

The reference – ultrathin CoFeB low current STT

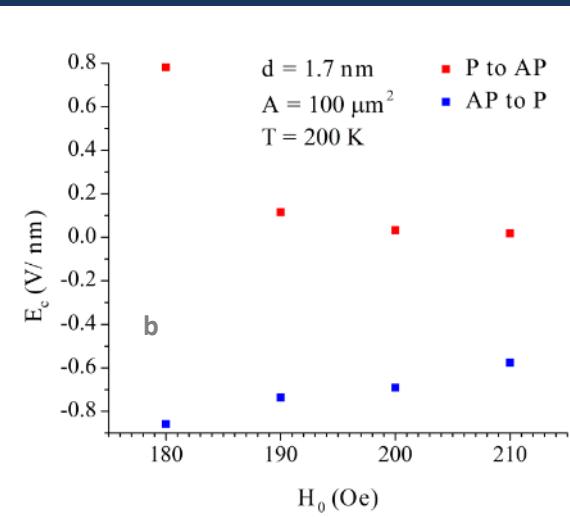
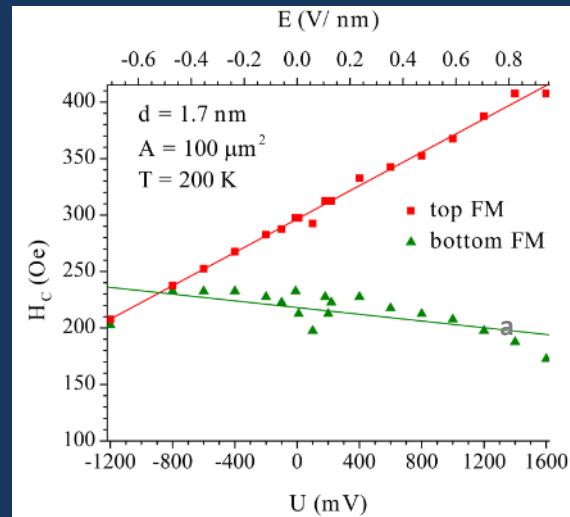
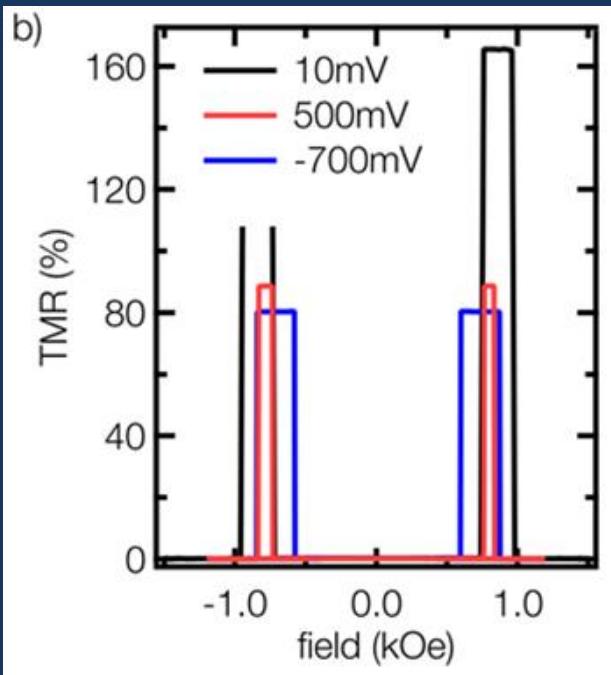
Critical current for STT switching the MTJ as a function of temperature (collaboration with M. Münzenberg group)



Temperature for
 $H_c \perp = 0$ from
last page
– 1nm layer has „no
anisotropy“

.. nearly zero critical current density possible by controlling temperature¹
- ? good for TAS-STT-MRAM²

The reference – ultrathin CoFeB E-field shift of H_c



- (a) Dependence of the coercive fields of the electrodes of an MTJ on the applied voltage for the top (1nm thick) and bottom electrode (1.2nm thick) and
- (b) Electric field leading to a magnetization switching at an applied bias field H_0 (A. Gebauer, Bachelor thesis, publication in preparation)

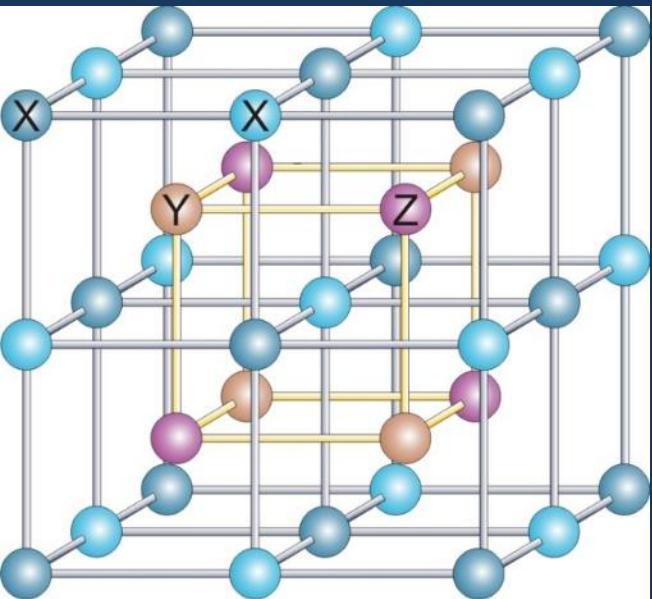
.. H_c shift about four times larger for 1.0nm than for 1.2nm CoFeB
.. Open: temperature dependence, possible synergy with STT ...

Why look for new Materials ?

But:

- „Ultrathin“ requires expensive equipment and process control
- The perpendicular magnetization comes from the interfaces between MgO and CoFeB (sensitive property)
- Thin CoFeB has a relatively large magnetic damping (speed issue !)
- Base layer (antiferromagnet) is expensive
- Narrow production window, relatively expensive, low speed
- Need new materials !





Heusler compounds (and related):

- X_2YZ composition
- crystallographic $L2_1$ structure
- high spin polarization / TMR ratios
- high Curie temperature T_C
- well known with in plane magnetization

N. Tezuka et al., Appl. Phys. Lett. **94**, 162504 (2009)

P. Webster, J. Phys. Chem. Solids **32**, 1221 (1971)

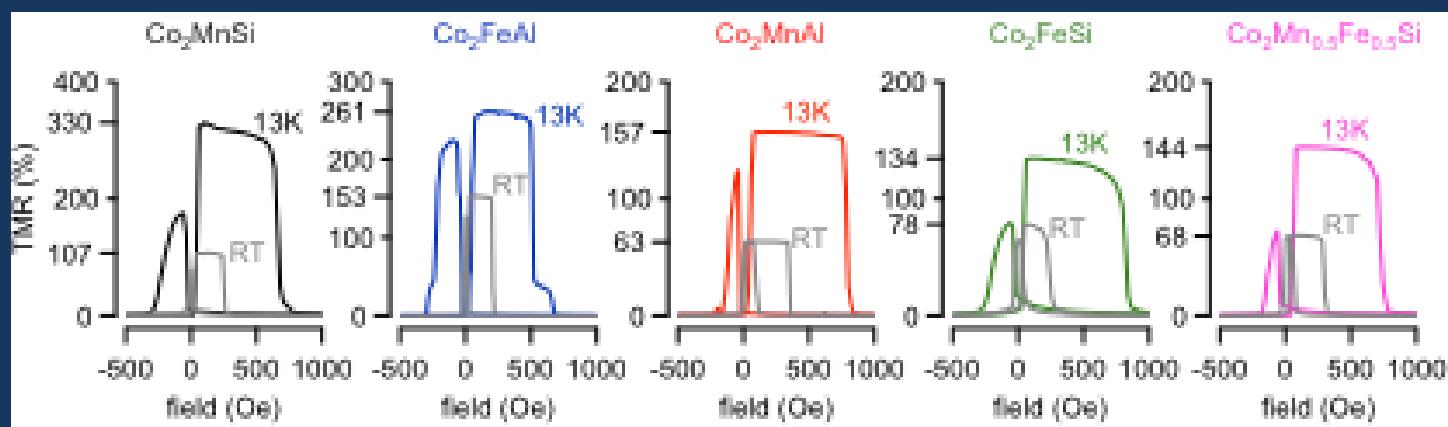
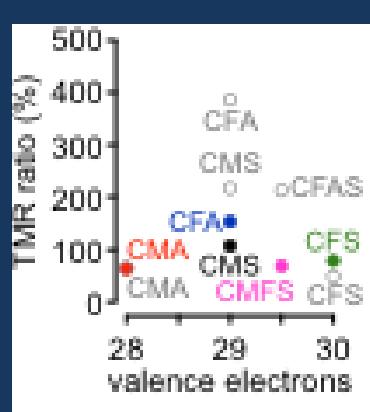
Bielefeld University
contributes > 20 "Heusler"
publications since 2006

D. Ebke, PhD thesis (2010)

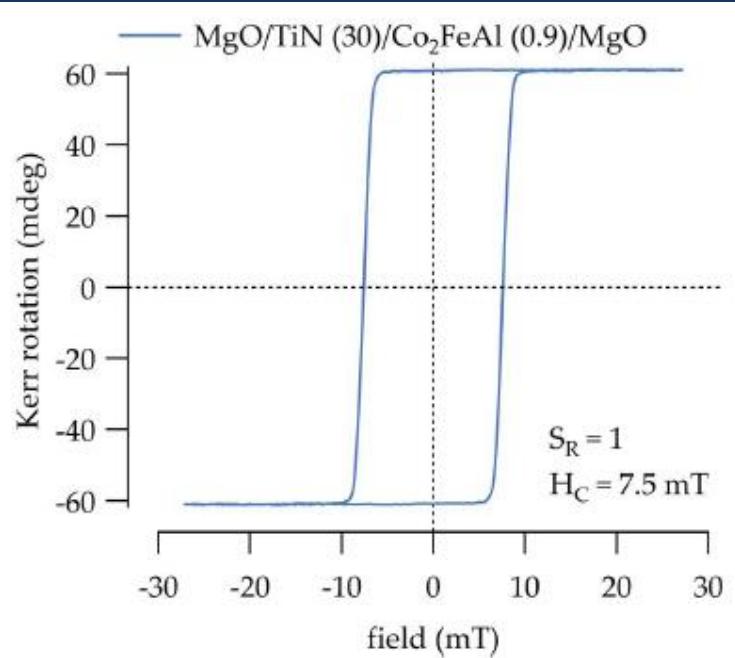
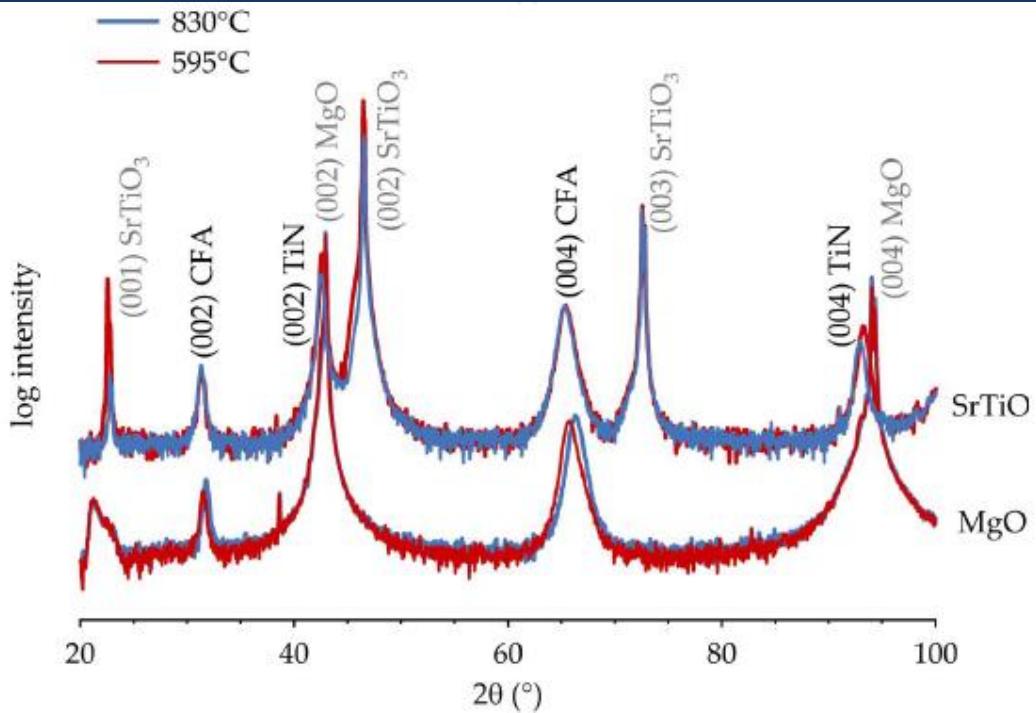
General goals:

- thermal stability ($KV > 50 - 60 \text{ k}_B T$)
- switching current low ($0.1-1 \text{ MA/cm}^2$)
- TMR ratio (100 - 200% @ RT, better more)
- fast (nsec or faster)
- cheap, reliable, easy to prepare, ...

Heusler electrodes in TMR cells with in plane anisotropy



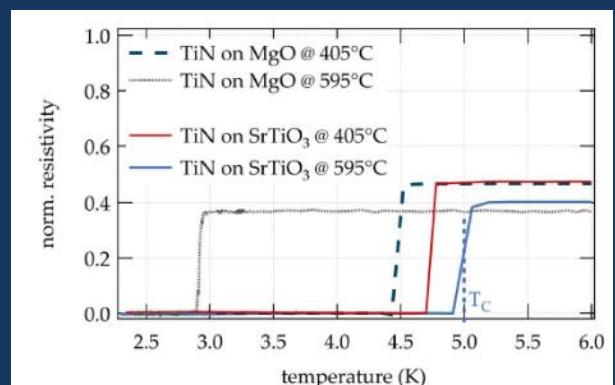
Very large TMR possible: H.X. Liu et.al., Appl. Phys. Lett. 101, 132418 (2012)
(2000% @ LT, 350% @ RT, $\text{Co}_2\text{Mn}_x\text{Si}$, $x=1.3$)



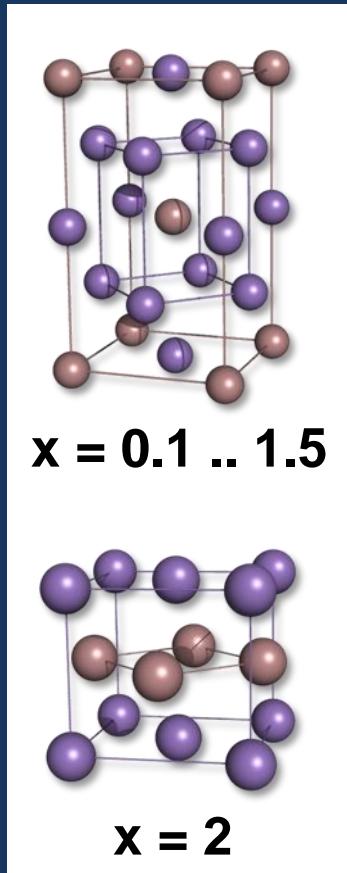
Perpendicular to plane

→ Ultrathin Co_2FeAl Heusler compound on TiN gives smooth surface and perpendicular anisotropy

→ TiN is a good conductor, very stable (and superconducting)



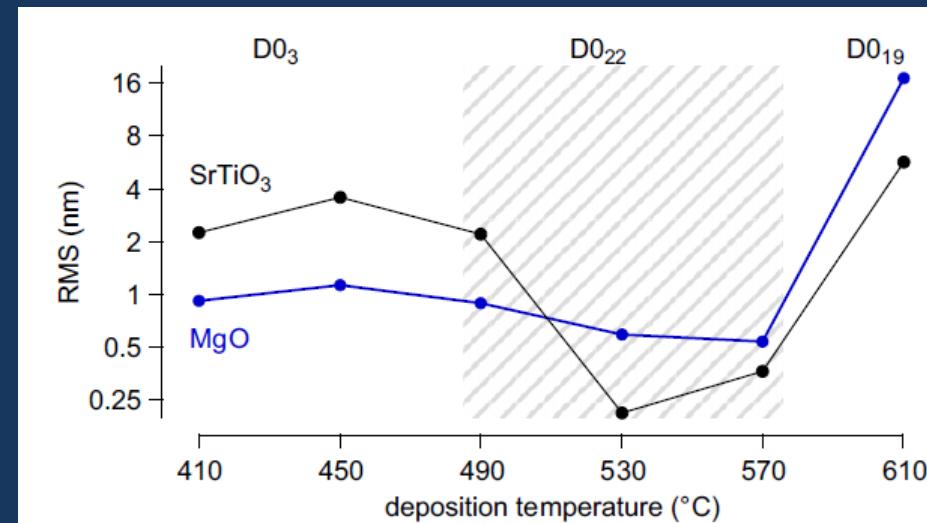
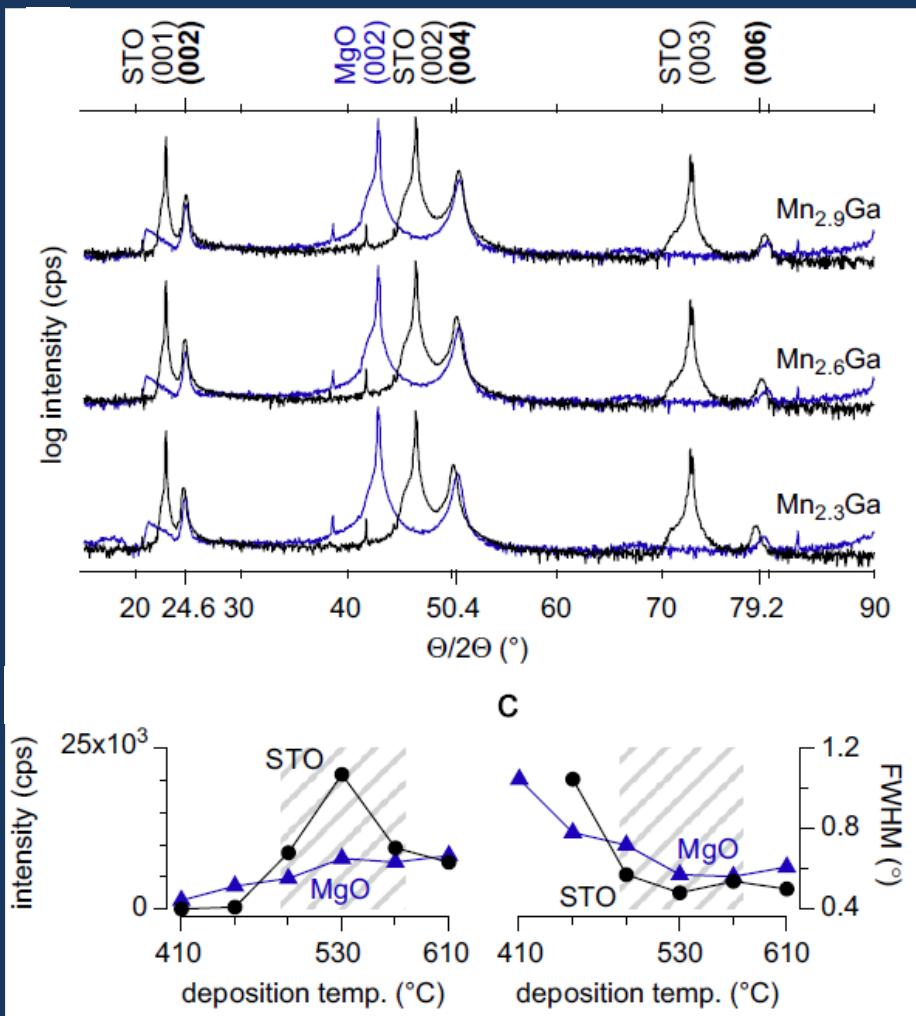
Heusler with perpendicular cryst. anisotropy: The Mn_xY family



Example Mn_{3-x}Ga

- predicted high spin polarization ($P=88\%$)
- perpendicular properties down to 5nm thickness
- high Curie temperature ($T_C \approx 770\text{K}$)
- low magnetic moment (about $0.26\mu\text{B} / \text{f.u.}$)
- tunable magnetic behavior:
 H_C decreases with increasing x (leak of Mn)

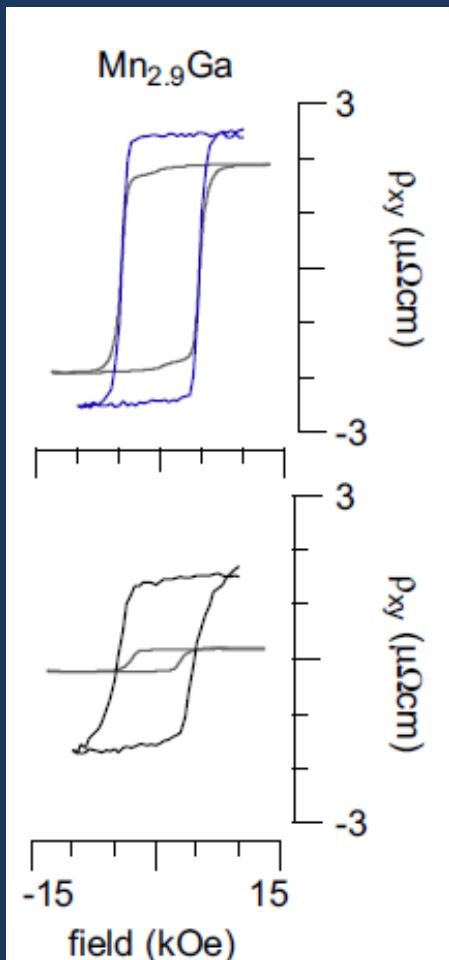
New Materials: Mn_xGa



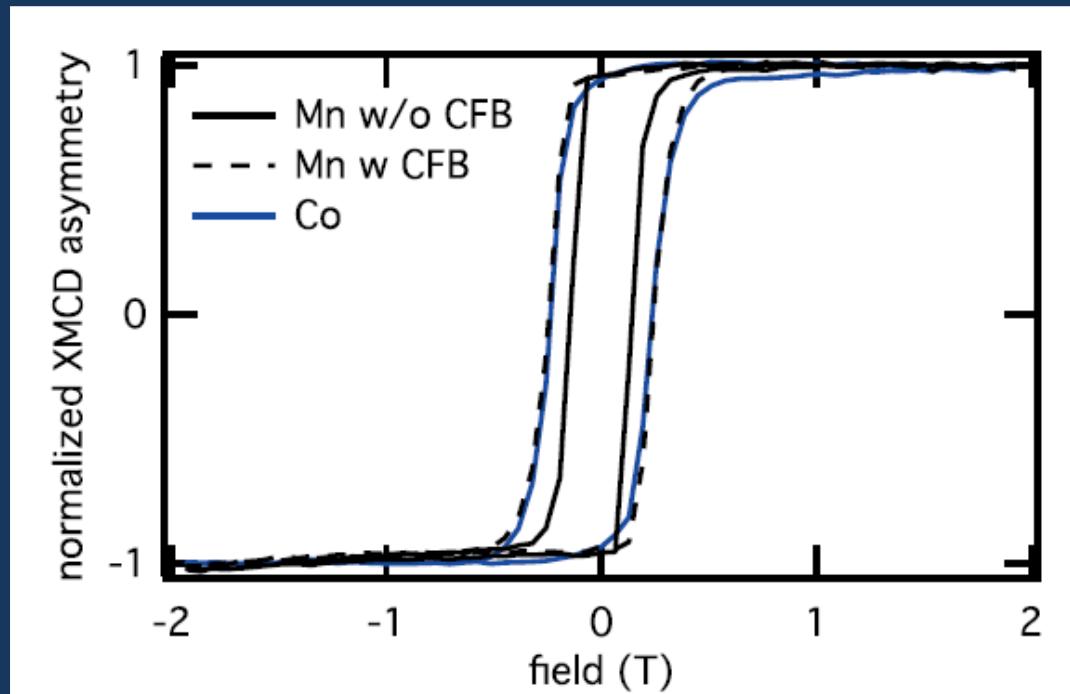
X-ray diffraction (left) of Mn_xGa

and

roughness (top) for Mn_{2.9}Ga

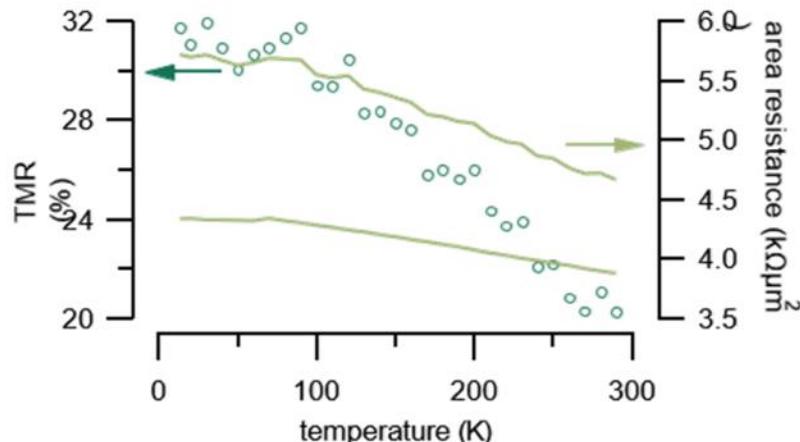
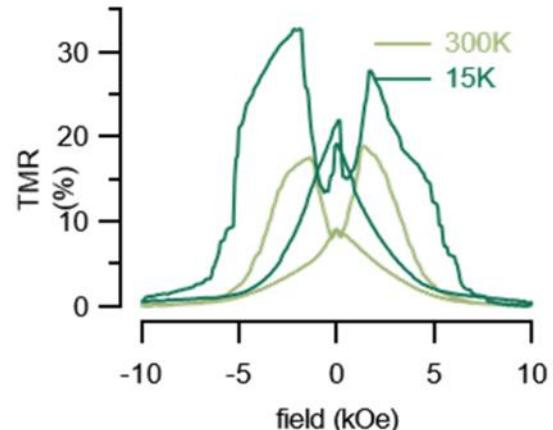
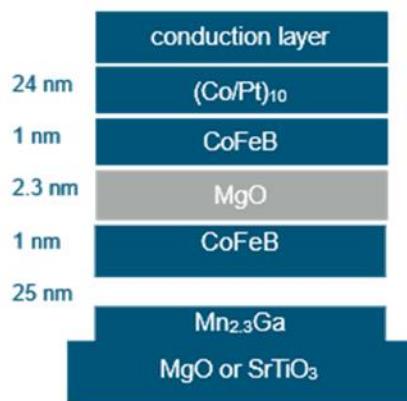


Overlay of magnetization (grey) and transversal Hall resistivity for $Mn_{2.9}Ga$ deposited on MgO (top) and STO (bottom) substrates at 530°C.



Element-specific hysteresis loops.
Dashed black → normalized Mn XMCD asymmetry for a sample with (without)
1 nm CoFeB @ RT, normal incidence, out-of-plane magnetic field.
Blue curve → normalized Co XMCD asymmetry.

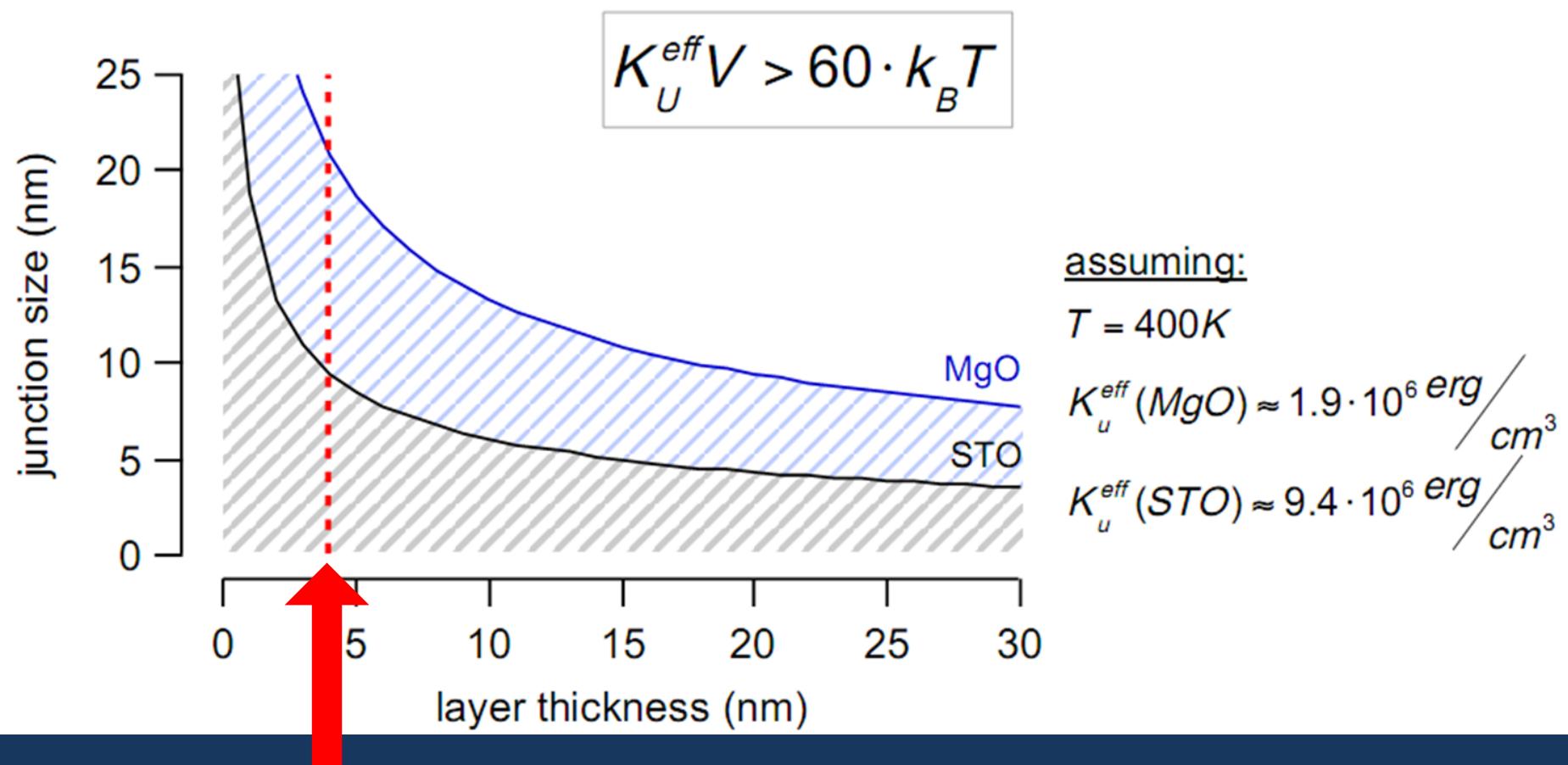
New Materials: Mn_xGa



- first reported TMR effect for perpendicular Mn-Ga compound¹
- highest TMR effect for samples with $\{\text{Co/Pt}\}_{10}$ multilayer counter electrode
- other groups reported higher TMR effects *for different interlayer*²
- no post-annealing process took place
- reasonable TMR ratio **only** for samples with ferromagnetic interlayer

[1] Glas, Integration of $Mn_{3-x}Ga$ Thin Films Into Magnetic Tunnel Junctions, Diploma Thesis, Bielefeld, 2012

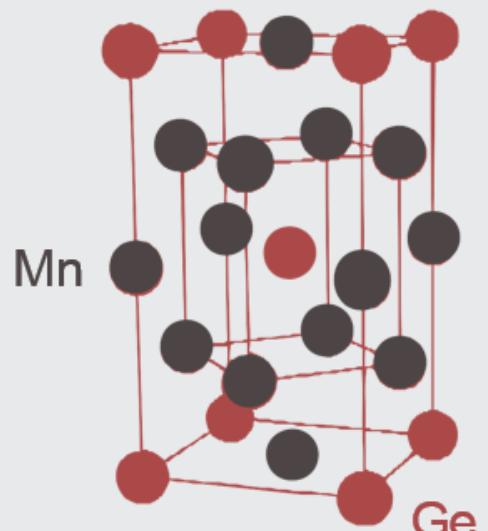
[2] Ma et al., J. Appl. Phys. 114, 163913 (2013).



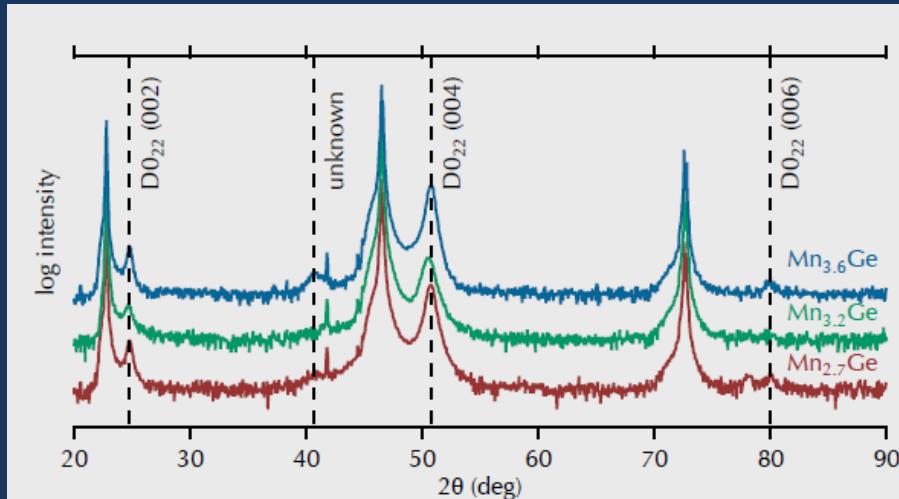
Could work down to 5nm feature size !

Mn_{3±x}Ge properties

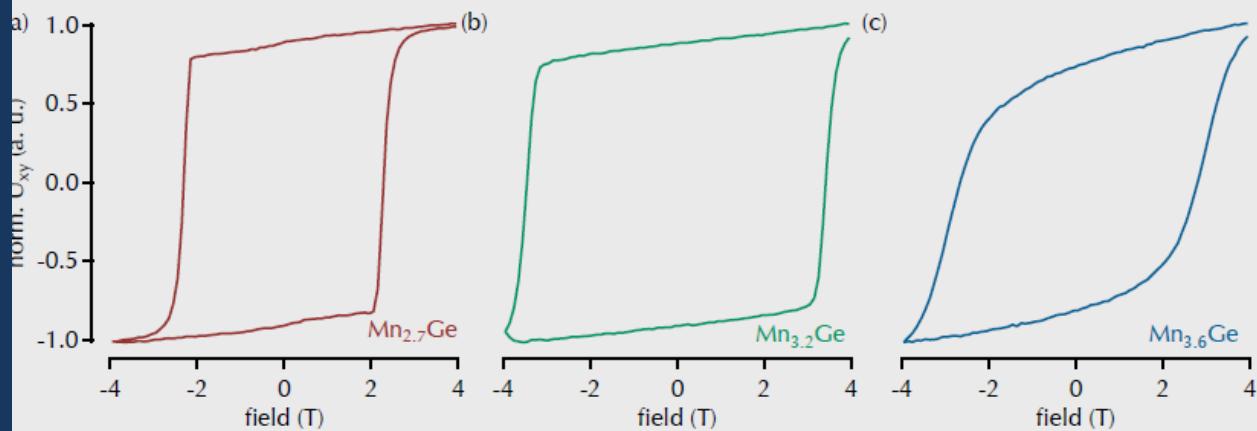
- Mn_{3±x}Ge compound crystallises in D0₂₂ phase
- lattice constants:
 - a,b = 3.816Å
 - c = 7.261Å
- magnetic moment:
 - 0.4μ_B/f.u. corresponds to 175kAm⁻¹
- Anisotropy constant between 0.9 and 1.2MJm⁻³ was reported
- spin polarisation of 46% via point contact Andreev reflection
- no single crystal phase was obtained by other groups

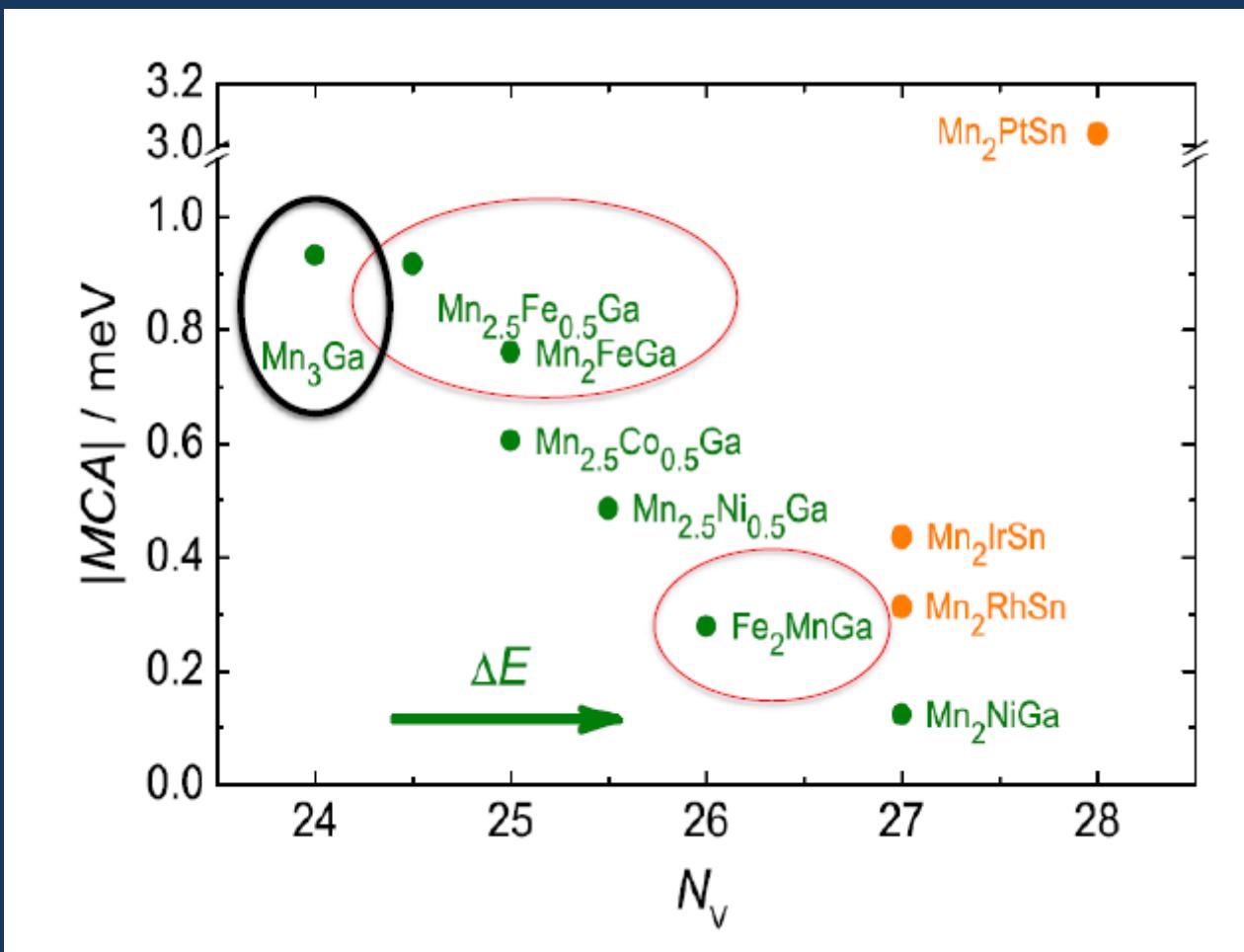


New Materials: Mn_xGe



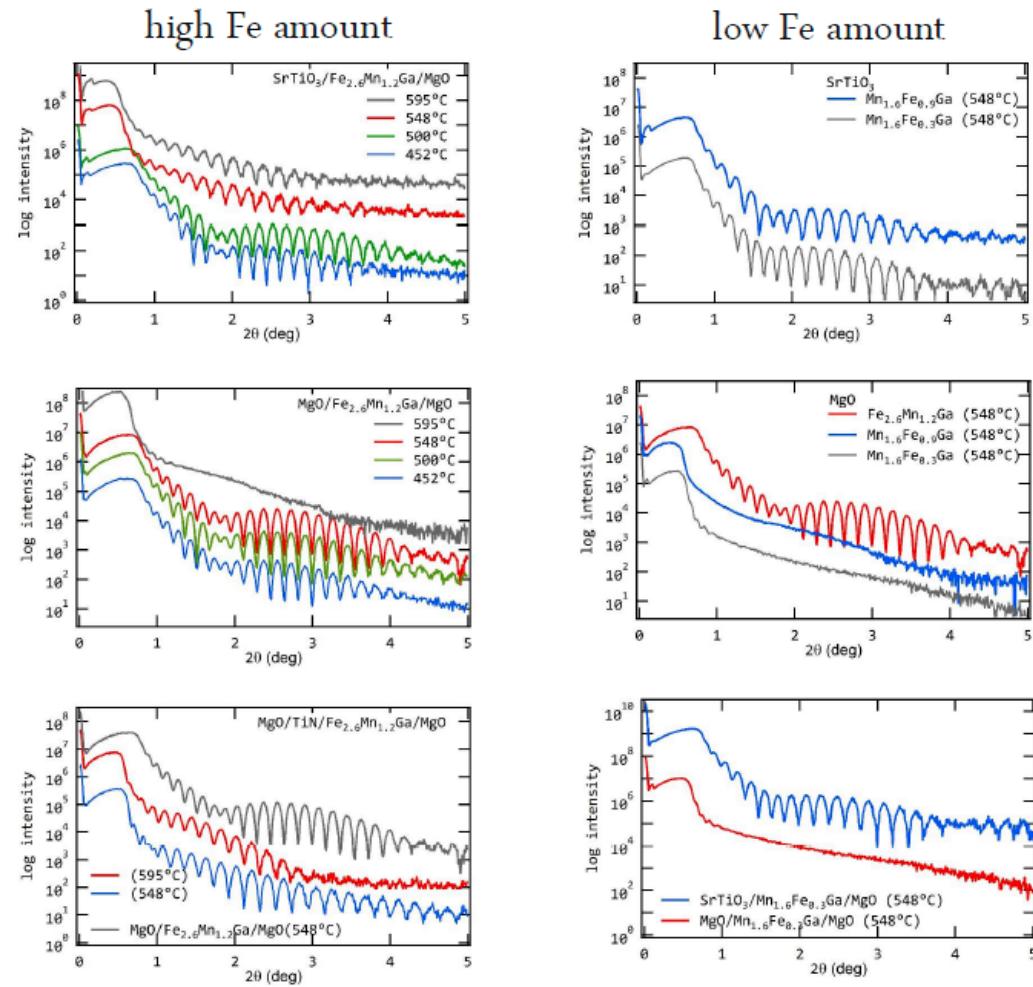
- D0₂₂ structure was achieved for different compositions
- Higher crystallinity for Mn rich samples
- Unknown reflex for the Mn_{3.6}Ge sample
- Highest H_c of 3.25T for Mn_{3.2}Ge
- upcoming in-plane component for Mn_{3.6}Ge
- critical Mn content achieved

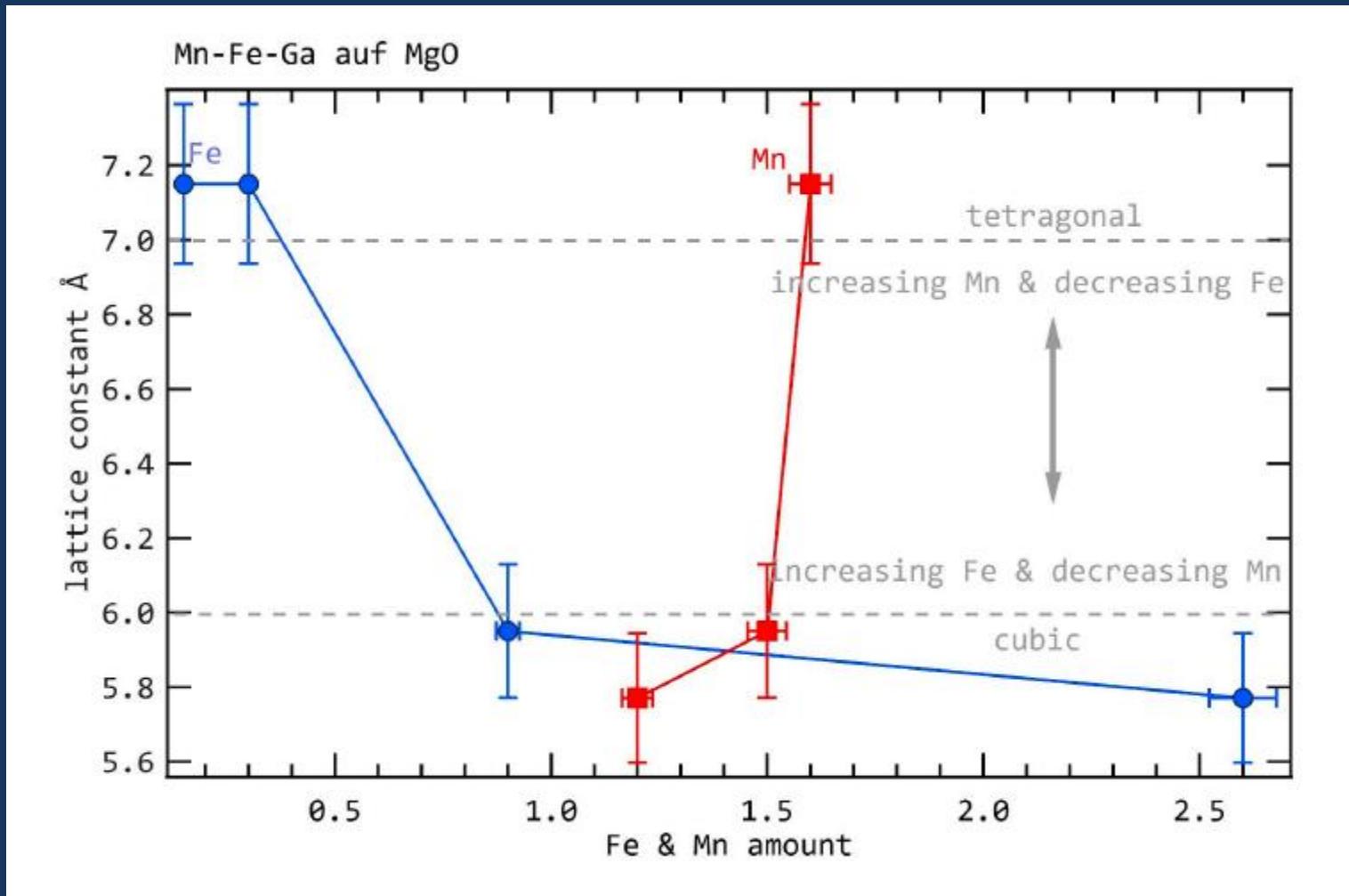




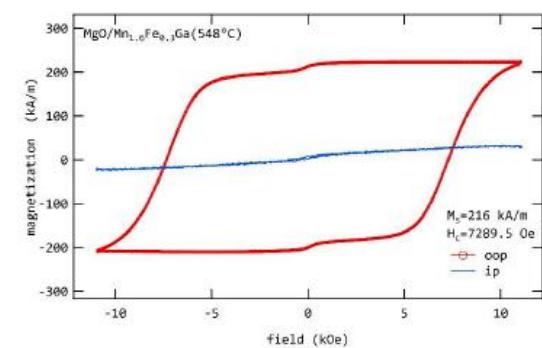
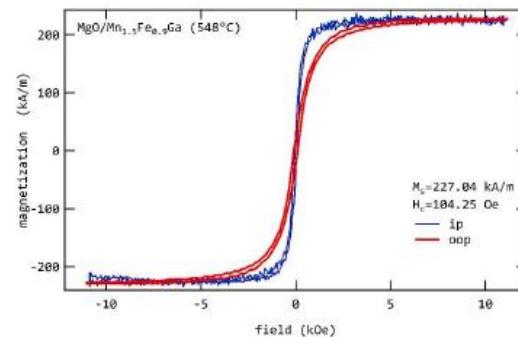
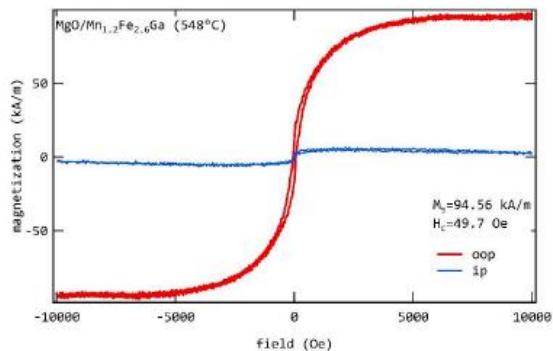
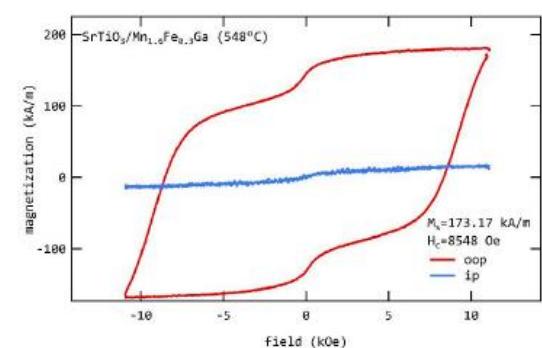
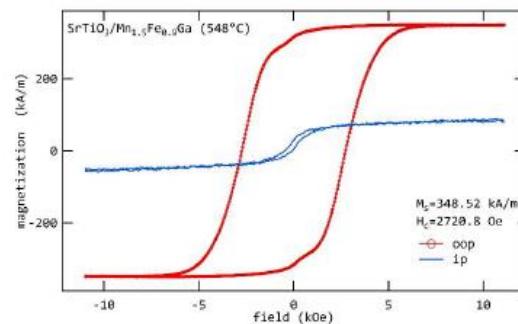
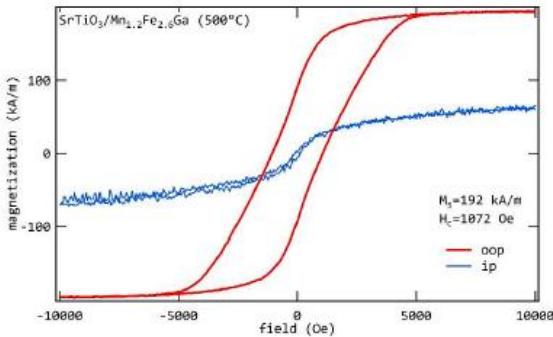
CRYSTALLINE PROPERTIES

- deposition time = 220 sec
- thickness: $d \approx 46 \pm 0.4 \text{ nm}$
- roughness (measured via XRR):
 $\approx 0.5 \pm 0.05 \text{ nm}$
 (high Fe amount)
- low Fe amount leads to high
 roughness (measured by AFM)





MAGNETIC PROPERTIES



- High H_c and oop-magnetization for deposition on SrTiO_3 and low Fe amounts on MgO

Markus Meinert group:

MnN

of the „ Mn_xY - family“ is anti-ferromagnetic and shows excellent exchange bias to CoFe:

A new exchange bias system !

MnN grown by reactive sputtering with Ar/ N_2 mixture

Can be grown on single crystalline substrates and on Si/SiO₂

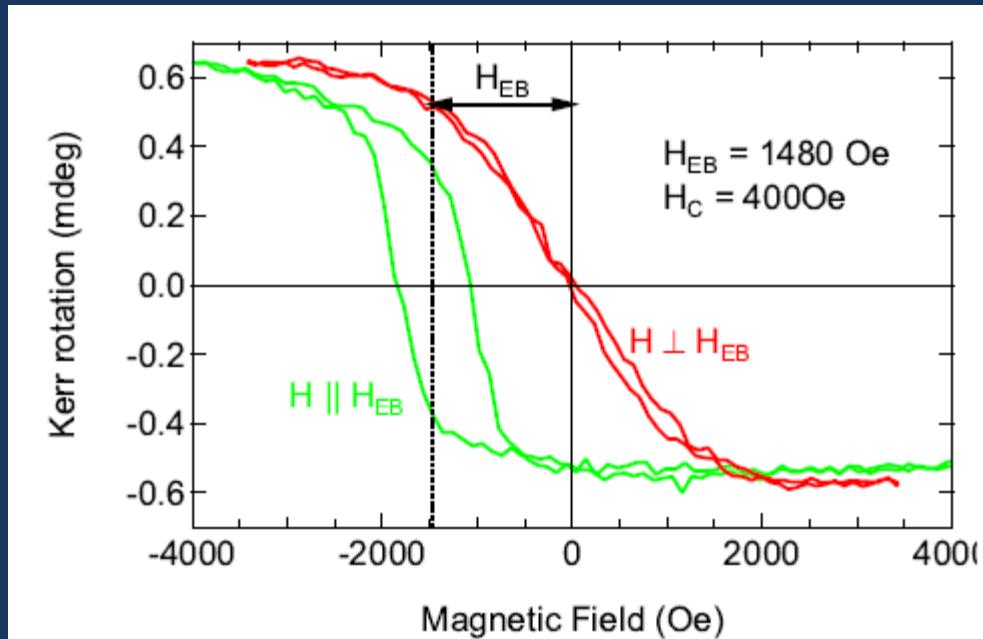
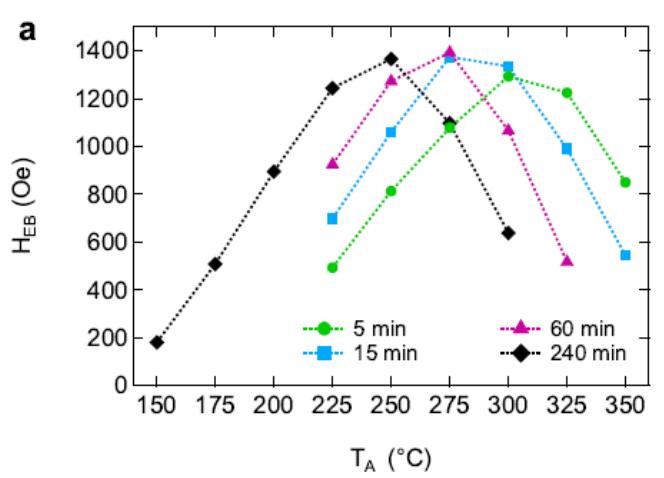
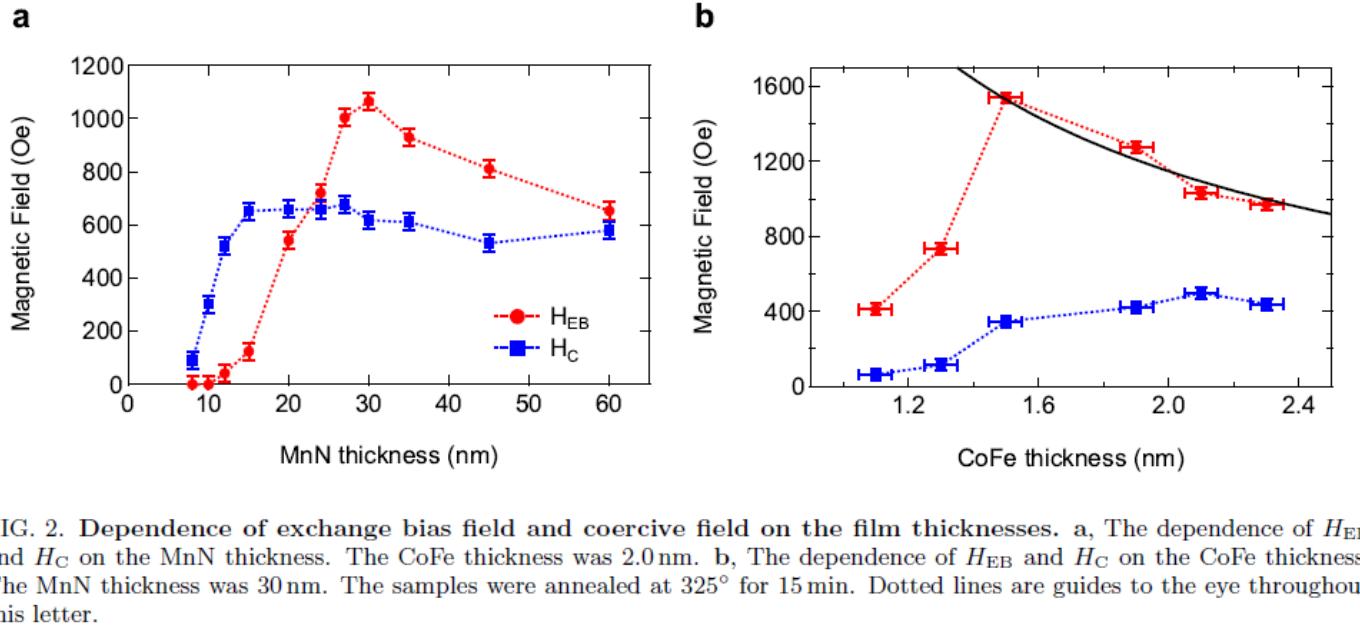


FIG. 1. Hysteresis loops of a MnN/CoFe stack at RT. The film thicknesses were $t_{\text{MnN}} = 30$ nm and $t_{\text{CoFe}} = 1.5$ nm. The sample was annealed at 325° for 15 min and field cooled. The definition of the exchange bias field H_{EB} is shown. With the external magnetic field perpendicular to the exchange biasing field no hysteresis shift is observed.

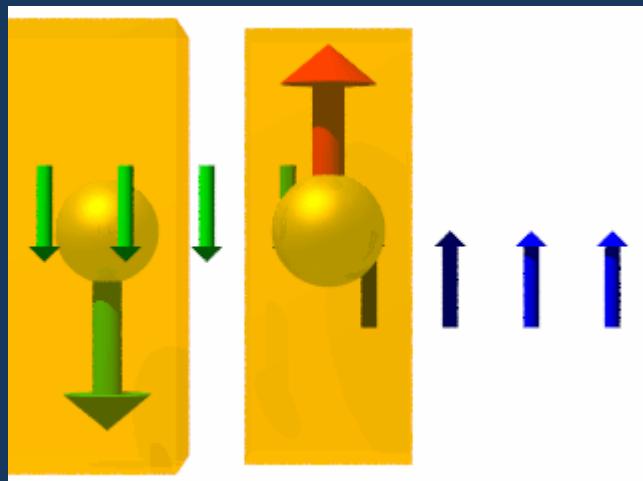


Dependence of the exchange bias field on annealing temperature and duration. The samples with MnN 30nm / CoFe 1.9nm were annealed and field cooled at temperature T_A for different times t_A .

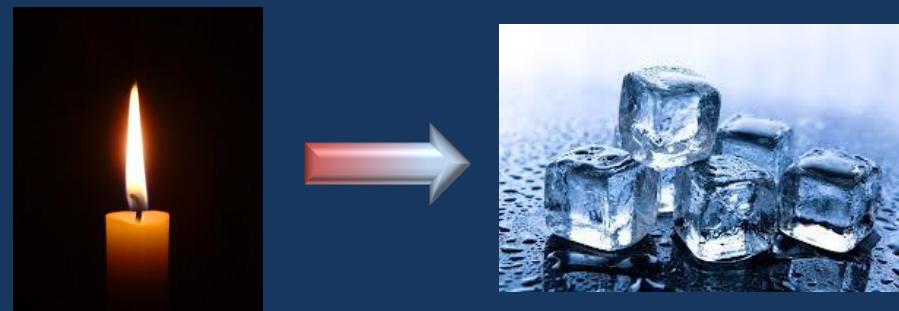
a, Samples successively heated for t_A with increasing temperature T_A .

b, Same data, parametrized with annealing temperature T_A .

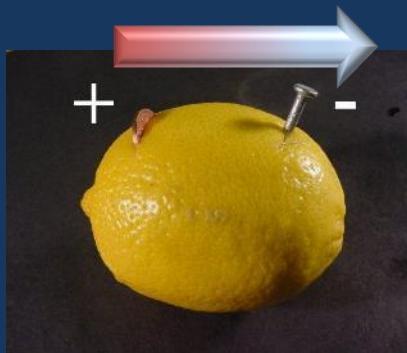
Remember: Spin Transfer Torque:



Can we drive Spin Currents also by other external forces?

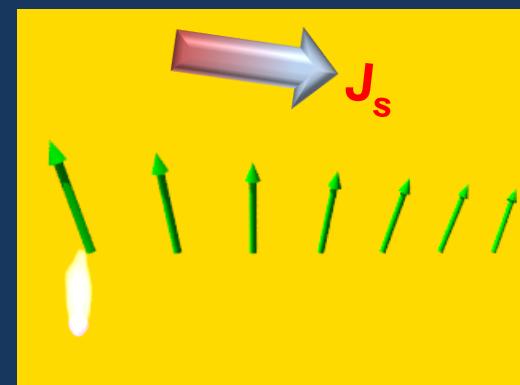


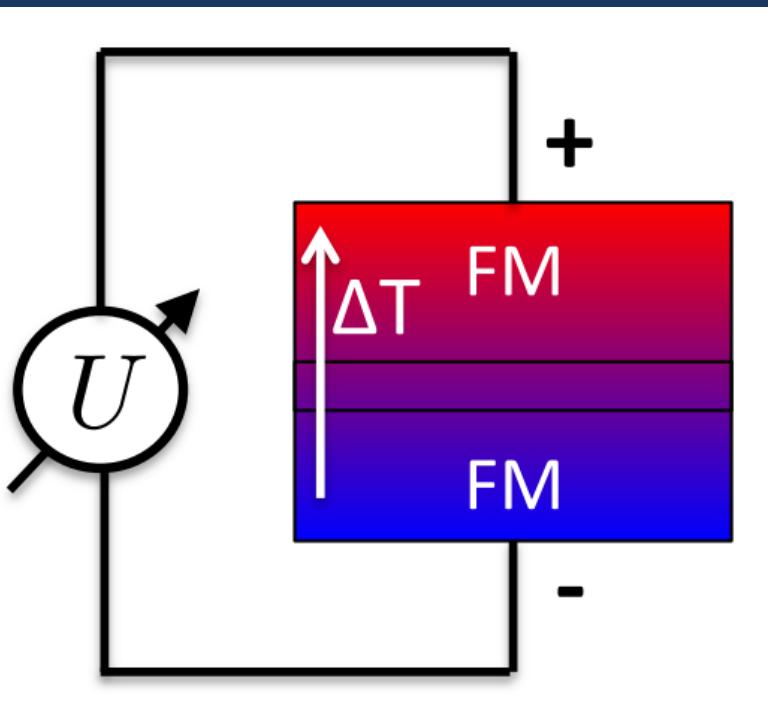
Spin Current is carried by electrons
Electrons are driven by voltage



$$\vec{E} = -\nabla\Phi(\vec{r})$$

YES: Temperature differences
 $\Delta T = \nabla T(\vec{r})$





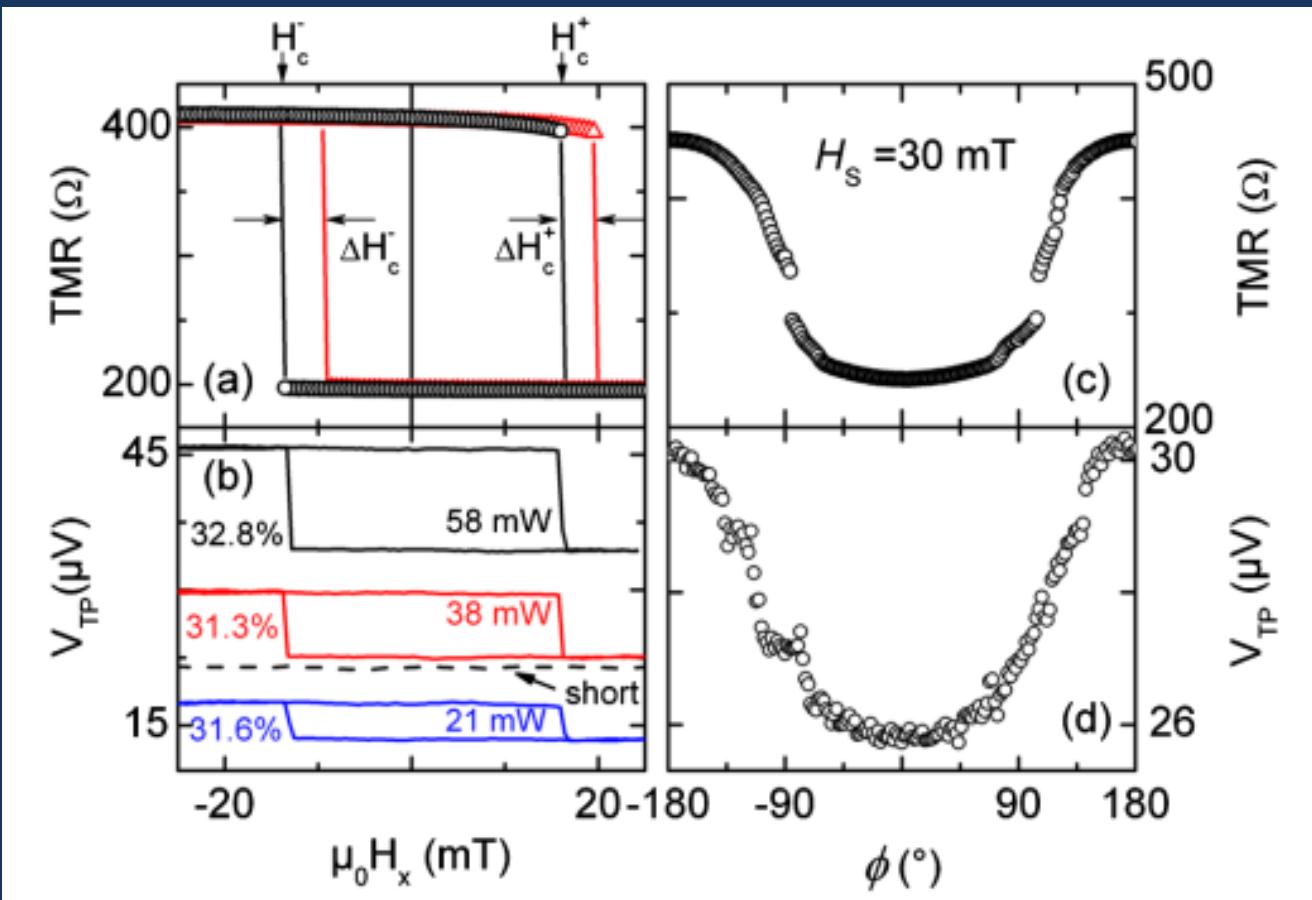
Seebeck coefficient for such tunnel devices:

$$S = \frac{\int T(E)(E - \mu)(-\partial_E f(E, \mu, T))dE}{e T \int T(E)(-\partial_E f(E, \mu, T))dE}$$

$\partial_E f(E, \mu, T)$: Derivative of occupation function

Thermovoltage should depend on magnetization directions

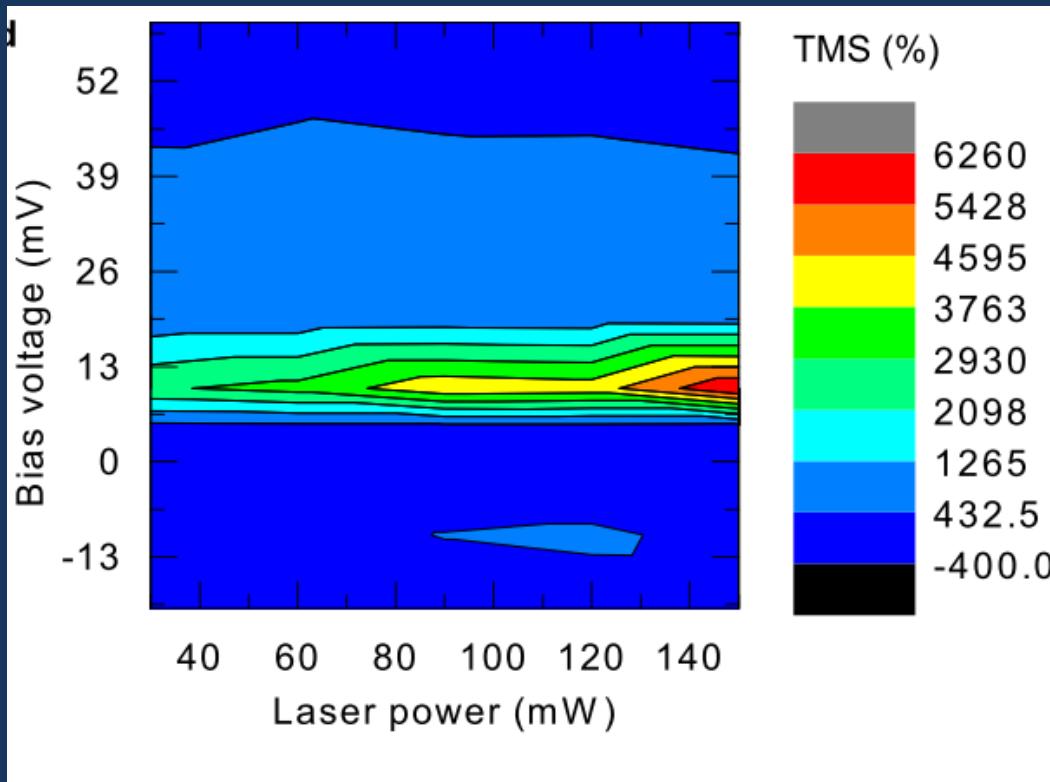
Important: S unequal to conductivity $g = \frac{e^2}{h} \int T(E)(-\partial_E f(E, \mu, T))dE$



M. Walter et.al., Nature Materials,
10 (2011) 742
(Münzenberg group Göttingen)

N. Liebing et.al., Phys. Rev. Lett.
107 (2011) 177201
(Schumacher group Braunschweig)

Tunneling
Magnetoresistance
and
Thermovoltage
of Magnetic
Tunnel Junctions



Additional bias voltage at
CoFeB/MgO/CoFeB
tunnel-junctions:
Tunnel Magneto
Seebeck effects
> 6000%
.. ongoing work

On/off switching of bit readout in bias-enhanced
tunnel magneto-Seebeck effect

Alexander Bohnke, Marius Milnikel, Marvin von der Ehe, Christian Franz, Vladyslav Zbarsky, Michael Czerner, Karsten Rott, Andy Thomas, Christian Heiliger, Günter Reiss & Markus Münzenberg

[Affiliations](#) | [Contributions](#) | [Corresponding author](#)

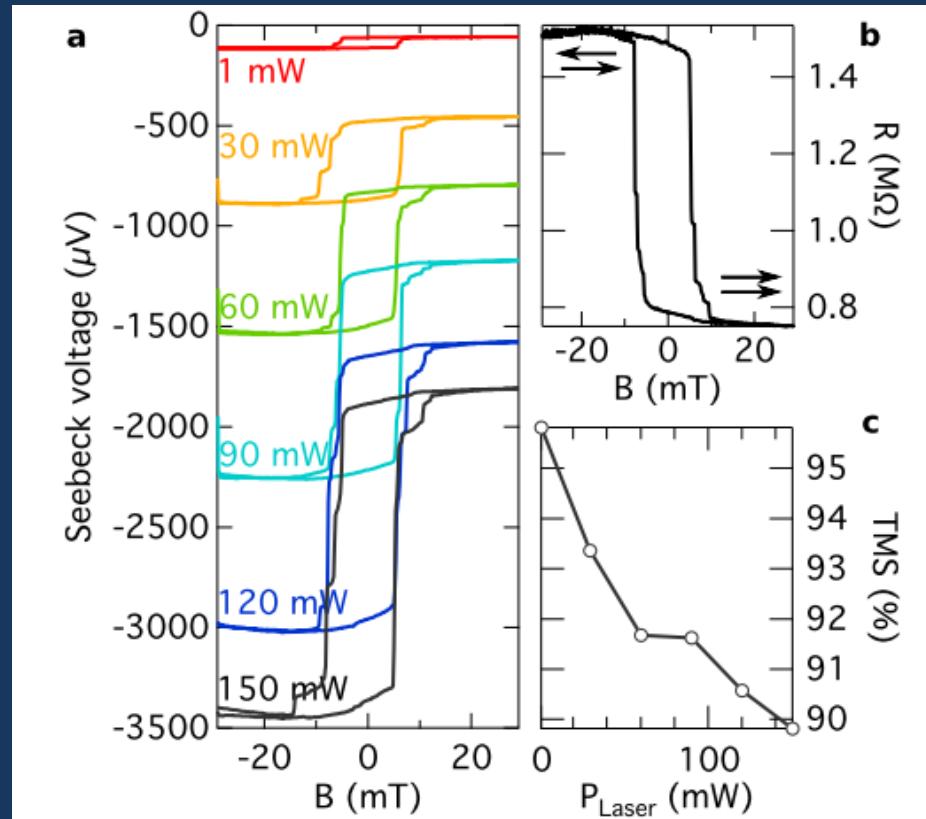


Gap in one spin direction should increase not only TMR but also

$$S = \frac{\int T(E)(E - \mu)(-\partial_E f(E, \mu, T))dE}{e T \int T(E)(-\partial_E f(E, \mu, T))dE}$$

(large asymmetry of DOS at E_F)

Spincalorics : MTJs + Heuslers



- a) TMS reaches 90 ... 96 % comparable to TMR (b)
- c) Dependence of TMS ratio on applied laser power.

... ongoing experimental and theoretical work



All coworkers in Bielefeld

Siemens AG, Sensitec, Prema,
Qiagen, Singulus, ...

M. Münzenberg, Göttingen
H.-W. Schumacher, PTB Braunschweig
S. Demokritov, Münster
H. Ebert, München
C. Felser, G. Jacob, G. Fecher, Mainz
B. Hillebrands, Kaiserslautern
J. Moodera, Cambridge
C. Chappert, Paris
R. Cowburn, London, ...

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